

EFFECT OF SORET AND TEMPERATURE DEPENDENT VISCOSITY ON THERMOHALINE CONVECTION IN A FERROFLUID SATURATING A POROUS MEDIUM

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Soret driven ferrothermoconvective instability in multi-component fluids has a wide range of applications in heat and mass transfer. This paper deals with the theoretical investigation of the effect of temperature dependent viscosity on a Soret driven ferrothermoconvective convection heated from below and salted from above subjected to a transverse uniform magnetic field in the presence of a porous medium. The Brinkman model is used in the study. It is found that the stationary mode of instability is preferred. For a horizontal fluid layer contained between two free boundaries an exact solution is examined using the normal mode technique for a linear stability analysis. The effect of salinity has been included in magnetization and density of the fluid. The critical thermal magnetic Rayleigh number for the onset of instability is obtained numerically for sufficiently large values of the buoyancy magnetization parameter M_1 using the method of numerical Galerkin technique. It is found that magnetization and permeability of the porous medium destabilize the system. The effect of temperature dependent viscosity stabilizes the system on the onset of convection.

Keywords: ferroconvection, porous medium, Soret effect, multi-component system, Brinkman model, temperature dependent viscosity, Galerkin technique.

1. Introduction

Magnetic fluids, also called ‘ferrofluids’, are electrically non-conducting colloidal suspensions of tiny particles of solid ferromagnetic material in a non-electrically conducting carrier fluid like water or heptanes, kerosene, hydrocarbon, etc. These fluids behave as a homogeneous continuum and exhibit a variety of interesting phenomena. Ferromagnetic fluids are not found in nature but are artificially synthesized. The viscosity of a magnetic nanofluid as a function of the applied magnetic field, direction of magnetic field with respect to the flow direction and temperature, is useful for endurable applications for magnetic inkjet printers, heat transfer, nanomotors, nanogenerators, inertial dampers, switches, sensors, transformer cooling, loudspeaker, similar micro-and nanofluidic devices, magnetic targeted drug delivery, cancer treatment in biomedicine field, etc. (Odenbach and Thurm, 2012), (Berkovsky and Bastovoy, 1996) and (Gazeau *et al.*, 1997).

An introduction to the research on magnetic fluids has been given in the monograph by Rosensweig (1985), which reviews several applications of heat transfer through ferrofluids, such as enhanced convective cooling having a temperature dependent magnetic moment due to magnetization of the fluid. This magnetization is called ferroconvection, which is similar to Bénard convection (Chandrasekhar, 1985). Convective instability of ferromagnetic fluids has been predicted by Finlayson (1970). Schwab *et al.* (1983) investigated experimentally the Finlayson’s problem in the case of a strong magnetic field and detected the onset of convection by plotting the Nusselt number versus the Rayleigh number. Then, the critical Rayleigh

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number corresponds to a discontinuity in the slope. Later, Stiles and Kagan (1990) examined the experimental problem reported by Schwab *et al.* (1983) and generalized Finlayson's model assuming that under a strong magnetic field, the rotational viscosity augments the shear viscosity. Furthermore, Vaidyanathan *et al.*, (1991) investigated the theoretically the convective instability of a ferromagnetic fluid in a porous medium of large permeability by use of the Brinkman model. This investigation has been done to the effect of temperature dependent viscosity by Ramanathan and Muchikel (2006) using Galerkin technique.

Suresh Govindan *et al.* (2012) made a numerical analysis on ferroconvection with temperature dependent viscosity and anisotropic porous medium. Nanjundappa *et al.* (2012) introduced the effect of temperature dependent viscosity on Marangoni-Bénard ferroconvection without a porous medium under microgravity conditions in a horizontal ferrofluid layer in the presence of a uniform vertical magnetic field. Moreover, this work has been analyzed to the effect magnetic field dependent viscosity in the absence of temperature dependent viscosity by Nanjundappa *et al.* (2010). They used the Rayleigh Ritz method with Chebyshev polynomials of second kind as trial function. The onset of buoyancy-driven convection in a ferromagnetic fluid in the presence of a porous medium was studied by Shivakumara *et al.* (2010). The thermorheological effect of magnetoconvection in fluids with weak electrical conductivity was studied numerically by Siddheshwar (2004).

The study of convection in two component ferrofluids will throw light on convective instability. This is referred to as a type of convection known as ferrothermohaline convection studied by Baines and Gill (1969). Vaidyanathan *et al.* (1995; 1997) investigated the presence and absence of a porous medium on ferrothermohaline. Here, they found that the salinity of a ferromagnetic fluid enabled the fluid to destabilize more when it is salted from above. A really interesting situation from both a geophysical and a mathematical viewpoint arises when the layer is simultaneously heated from below and salted from above. Sunil *et al.* (2011) considered the double-diffusive convection on ferromagnetic fluid with rotation and internal angular momentum.

Recently, Vaidyanathan *et al.* (2005) attempted to study the Soret effect due to ferrothermohaline convection of a sparse distribution. Further, the condition of a porous medium of ferroconvective instability of multi-component fluid heated from below and salted from above was analyzed by Sekar *et al.* (2006; 2013) for isotropic and anisotropic models. The effect of rotation on thermohaline convection in a ferromagnetic fluid saturating an anisotropic porous medium with the Soret effect was obtained by Sekar *et al.* (2013a) and this investigation was carried out for magnetic field dependent viscosity by Sekar and Raju (2013).

In view of these investigations, it is attempted to analyze the effect of temperature dependent viscosity on the Soret driven ferrothermohaline convection in the presence of an isotropic porous medium of low permeability, subjected to a vertical magnetic field using the Brinkman model. In this investigation, the free boundaries are considered. The resulting eigen value problem is solved numerically using the Galerkin method. Besides, an analytical formula is obtained for the critical magnetic Rayleigh number by a regular perturbation method.

2. Mathematical formulation

In this investigation, we consider an infinite spread horizontal layer of an Oberbeck-Boussinesq ferromagnetic fluid of thickness " d " saturating a sparsely distributed porous medium heated from below and salted from above. The temperature and salinity at the bottom and top surfaces are $z = \pm d/2$ are $T_0 \pm \Delta T/2$ and $S_0 \pm \Delta S/2$, respectively. Both the boundaries are assumed to be free and perfect conductors of heat and salt. This fluid layer is taken to be an isotropic porous medium and the fluid viscosity is assumed to be temperature-dependent in the following form (Ramanathan and Muchikel, 2006 and Siddheshwar, 2004)

$$\mu(T) = \mu_I \left[1 - \delta (T - T_a)^2 \right] \quad (2.1)$$

where δ is a small positive quantity.

The gravity field $\mathbf{g} = (0, 0, -g)$ and uniform vertical magnetic field intensity $\mathbf{H} = (0, 0, H_0)$ pervade the system. Considering the Soret effect on the temperature gradient the mathematical equations governing the above investigation are as follows.

The continuity equation for an incompressible fluid is

$$\nabla \cdot \mathbf{q} = 0. \quad (2.2)$$

The corresponding momentum equation is

$$\rho_0 \frac{D\mathbf{q}}{Dt} = -\nabla p + \rho \mathbf{g} + \nabla \cdot (\mathbf{H}\mathbf{B}) + \nabla \left[\mu(T) (\nabla \mathbf{q} + \nabla \mathbf{q}^{Tr}) \right] - \frac{\mu(T)}{k} \mathbf{q}. \quad (2.3)$$

The temperature equation for an incompressible ferromagnetic fluid is

$$\left[\rho_0 C_{v,H} - \mu_0 \mathbf{H} (\partial \mathbf{M} / \partial T)_{v,H} \right] (dT/dt) + \mu_0 T (\partial \mathbf{M} / \partial T)_{v,H} \cdot (d\mathbf{H}/dt) = K_I \nabla^2 T + \phi. \quad (2.4)$$

The conservation of mass flux equation is given by

$$\rho_0 (\partial / \partial t + \mathbf{q}) S = K_S \nabla^2 S + S_T \nabla^2 T. \quad (2.5)$$

The density equation of state for a Boussinesq two-component fluid is

$$\rho = \rho_0 \left[1 - \alpha_t (T - T_0) + \alpha_s (S - S_0) \right]. \quad (2.6)$$

Maxwell's equations, simplified for a non-conducting fluid with no displacement currents, become

$$\nabla \cdot \mathbf{B} = 0, \quad \nabla \times \mathbf{H} = 0 \quad (2.7a,b)$$

where the magnetic induction is given by

$$\mathbf{B} = \mu_0 (\mathbf{M} + \mathbf{H}). \quad (2.8)$$

In general, the pressure of ferromagnetic fluid can distort an external magnetic field if a magnetic interaction (dipole-dipole) takes place, but this is negligible for small particle concentration, as is assumed here. We assume that the magnetization is aligned with the magnetic field, but allows a dependence on the magnitude of the magnetic field, temperature and salinity, so that

$$\mathbf{M} = \frac{\mathbf{H}}{H} M(H, T, S). \quad (2.9)$$

The magnetic equation of state is linearized about the magnetic field H_0 , an average temperature T_0 and an average salinity S_0 , to become

$$M = M_0 + \chi(H - H_0) - K(T - T_0) + K_2(S - S_0). \quad (2.10)$$

Here H_0 is the uniform magnetic field of the fluid layer when placed in an external magnetic field $\mathbf{H} = H_0^{ext} \mathbf{k}$, \mathbf{k} is a unit vector in the z -direction, $H = |\mathbf{H}|$, $M = |\mathbf{M}|$ and $M_0 = (H_0, T_0, S_0)$.

The basic state is assumed to be quiescent and is given by

$$\mathbf{q} = \mathbf{q}_b = (0, 0, 0), \quad T = T_b = T_0 - \beta_t z, \quad S = S_b = S_0 - \beta_s z, \quad \rho(z) = \rho_0 [1 + \alpha_t \beta_t z - \alpha_s \beta_s z], \quad (2.11)$$

$$p = p_b(z), \quad H_b(z) = \left[H_0 - \frac{K\beta_t z}{1+\chi} + \frac{K_2\beta_s z}{1+\chi} \right] \mathbf{k}, \quad M_b(z) = \left[M_0 + \frac{K\beta_t z}{1+\chi} - \frac{K_2\beta_s z}{1+\chi} \right] \mathbf{k}$$

where \mathbf{k} is the unit vector in the vertical direction, β_t and β_s are non-negative constants.

Moreover, the basic state is disturbed by an infinitesimal thermal perturbation. Let the component of the perturbed magnetization and the magnetic field be $(M'_1, M'_2, M_b(z) + M'_3)$ and $(H'_1, H'_2, H_b(z) + H'_3)$, respectively. The perturbed viscosity and temperature are taken as $\mu_b(z) + \mu'$ and $T_b(z) + T'$, respectively. On linearization, and assuming $K\beta_t d \ll (1+\chi)H_0$ and $K_2\beta_s d \ll (1+\chi)H_0$ and using the expressions for H_b and M_b in Eqs (2.11), Eqs (2.7)-(2.9) become

$$H'_i + M'_i = [1 + (M_0/H_0)] H'_i, \quad \text{for } i = 1, 2, \quad (2.12)$$

$$H'_3 + M'_3 = (1 + \chi) H'_3 - K_2 S + S_T K \theta - K \theta,$$

$$B'_i = \mu_0 [1 + (M_0/H_0)] H'_i, \quad \text{for } i = 1, 2, \quad (2.13)$$

$$B_3 = \mu_0 (M_0 + H_0) + \mu_0 (1 + \chi) H'_3 - \mu_0 K_2 S + \mu_0 S_T K \theta - \mu_0 K \theta.$$

Equation (2.7b) implies that $\mathbf{H}' = \nabla \phi'$, where ϕ' is the perturbed magnetic potential and using the analyses of Sekar *et al.*, (2013; 2013a), the vertical component of the momentum equation can be written as

$$\begin{aligned} \rho_0 \frac{\partial}{\partial t} (\nabla^2 w) = & \rho_0 g \alpha_t \nabla_I^2 T' - \rho_0 g \alpha_s \nabla_I^2 S' - K \beta_t \frac{\partial}{\partial z} (\nabla_I^2 \phi') + \frac{\mu_0 K^2 \beta_t (1 - S_T)}{1 + \chi} \nabla_I^2 T' + \\ & - \frac{\mu_0 K K_2 \beta_s (1 - S_T)}{1 + \chi} \nabla_I^2 T' - \frac{\mu_0 K K_2 \beta_t}{1 + \chi} \nabla_I^2 S' + \frac{\mu_0 K^2 \beta_s}{1 + \chi} \nabla_I^2 S' + \mu_0 (\nabla^2 (\nabla^2 w)) + \\ & - \mu_0 K_2 \beta_s \frac{\partial}{\partial z} (\nabla_I^2 \phi') - \frac{\mu_b}{k} \nabla^2 w + \frac{\partial^2 \mu_b}{\partial z^2} \nabla^2 w + 2 \frac{\partial \mu_b}{\partial z} \frac{\partial}{\partial z} (\nabla^2 w) - \frac{1}{k} \frac{\partial \mu_b}{\partial z} \left(\frac{\partial w}{\partial z} \right). \end{aligned} \quad (2.14)$$

The modified Fourier heat conduction equation is

$$\rho_0 C_{v,H} \frac{\partial \theta}{\partial t} - \mu_0 K T_0 \frac{\partial}{\partial t} \left(\frac{\partial \phi}{\partial z} \right) = K_I (\nabla^2 \theta) + \left[\rho_0 c \beta_t - \left(\frac{\mu_0 K^2 T_0^2 \beta_t}{1 + \chi} \right) + \left(\frac{\mu_0 K K_2 T_0 \beta_s}{1 + \chi} \right) \right] w \quad (2.15)$$

where $\rho_0 C = \rho_0 C_{v,H} + \rho_0 K H_0$.

The salinity equation is

$$\frac{\partial S}{\partial t} + \beta_S w = K_S (\nabla^2 S) + S_T (\nabla^2 \theta). \quad (2.16)$$

Using the analysis similar to Sekar *et al.* (2013) one gets

$$(1 + \chi) \frac{\partial^2 \phi}{\partial z^2} + \left(1 + \frac{M_0}{H_0}\right) \nabla_I^2 \phi - K \frac{\partial \theta}{\partial z} + K_2 \frac{\partial S}{\partial z} + S_T K \frac{\partial \theta}{\partial z} = 0 \quad (2.17)$$

where $\nabla_I^2 = \left(\frac{\partial^2}{\partial x^2}\right) + \left(\frac{\partial^2}{\partial y^2}\right)$ and $\nabla^2 = \nabla_I^2 + \left(\frac{\partial^2}{\partial z^2}\right)$.

3. Normal mode analysis method

Analyzing the small thermal disturbances into normal modes, we assume that the perturbation quantities are of the form

$$(w, T, \phi, S) = [w(z, t), T(z, t), \phi'(z, t), S'(z, t)] \exp(i(k_x x + k_y y))$$

where $k_0 = \sqrt{k_x^2 + k_y^2}$ is the resultant wave number, k_x is the wave number along the x direction and k_y is the wave number along the y direction.

Following the normal mode analysis, the linearized perturbation dimensionless equations

$$\begin{aligned} \frac{\partial}{\partial t^*} (D^2 - a^2) w^* &= a R^{1/2} [M_1 D \phi^* - (1 + M_1 (1 - S_T) T^*)] + M_1 M_5 a R^{1/2} D \phi^* + \\ &- M_1 M_5 a R^{1/2} (1 - S_T) T^* + (D^2 - a^2)^2 w^* + a R_S^{1/2} [1 + M_4 + M_4 M_5^{-1}] S^* + \\ &+ (1 - Vz^{*2}) (D^2 - a^2)^2 w^* - \frac{(1 - Vz^{*2})}{k^*} (D^2 - a^2) w^* - 2V (D^2 - a^2) w^* + \\ &- 4Vz^* (D^2 - a^2) D w^* + \frac{2Vz^*}{k^*} D w^*, \end{aligned} \quad (3.1)$$

$$P_r \left[\frac{\partial T^*}{\partial t^*} - M_2 \frac{\partial}{\partial t^*} (D \phi^*) \right] = (D^2 - a^2) T^* + a R^{1/2} (1 - M_2 - M_2 M_5) w^*, \quad (3.2)$$

$$P_r \frac{\partial S^*}{\partial t^*} = \tau (D^2 - a^2) S^* - a R_S^{1/2} M_6 w^* + S_T M_5 M_6^{-1} (R / R_S)^{1/2} (D^2 - a^2) T^*, \quad (3.3)$$

$$D^2 \phi^* - M_3 a^2 \phi^* - (1 - S_T) D T^* + M_5 M_6^{-1} (R / R_S)^{1/2} D S^* = 0 \quad (3.4)$$

where the following non-dimensional parameters are introduced

$$\begin{aligned}
w^* &= \frac{wd}{v}, & t^* &= \frac{vt}{d^2}, & T^* &= \left(\frac{K_I a R^{1/2}}{\rho_0 C_{v,H} \beta_t v d} \right) \theta, & \phi^* &= \left(\frac{(1+\chi) K_I a R^{1/2}}{\rho_0 C_{v,H} K \beta_t v d^2} \right) \phi, \\
z^* &= \frac{z}{d}, & a &= k_0 d, & D &= \frac{\partial}{\partial z^*}, & S^* &= \left(\frac{K_S a R_S^{1/2}}{\rho_0 C_{v,H} \beta_S v d} \right) S, & v &= \frac{\mu}{\rho_0}, & k^* &= \frac{k}{d^2}, \\
M_1 &= \frac{\mu_0 K^2 \beta_t}{(1+\chi) \rho_0 g \alpha_t}, & M_2 &= \frac{\mu_0 K^2 T}{(1+\chi) \rho_0 C_{v,H}}, & M_3 &= \frac{I + (M_0/H_0)}{(1+\chi)}, & & & & & (3.5) \\
M_4 &= \frac{\mu_0 K^2 \beta_S}{(1+\chi) \rho_0 g \alpha_S}, & M_5 &= \frac{K_2 \beta_S}{K \beta_t}, & M_6 &= \frac{K_S}{K_I}, & P_r &= \frac{\mu C_{v,H}}{K_I}, \\
R &= \frac{\rho_0 C_{v,H} \beta_t \alpha_t g d^4}{v K_I}, & R_S &= \frac{\rho_0 C_{v,H} \beta_S \alpha_S g d^4}{v K_S}, & \tau &= \rho_0 C_{v,H} \left(\frac{K_S}{K_I} \right)
\end{aligned}$$

where R is the thermal Rayleigh number, R_S is the salinity Rayleigh number, Pr is the Prandtl number.

4. Exact solution for free boundaries using Galerkin technique

Here the simplest boundary conditions chosen, namely free-free, isothermal with infinite magnetic susceptibility χ in the perturbed field keep the problem analytically tractable and serve the purpose of providing a qualitative insight into the problem. The case of two free boundaries is of little physical interest, but it is mathematically important because one can derive an exact solution, whose properties guide our analysis. Thus the exact solution of the system subjected to the boundary conditions

$$w^* = D^2 w^* = T^* = D\phi^* = S^* = 0 \quad \text{at} \quad z^* = -1/2 \quad \text{and} \quad z^* = +1/2. \quad (4.1)$$

is written in the form

$$w^* = A w_1(z) e^{\sigma t^*} \cos \pi z^*, \quad T^* = B T_1(z) e^{\sigma t^*} \cos \pi z^*, \quad S^* = C S_1(z) e^{\sigma t^*} \cos \pi z^*, \quad (4.2)$$

$$D\phi^* = F \phi_1(z) e^{\sigma t^*} \cos \pi z^*, \quad \phi^* = \frac{F}{\pi} \phi_1(z) e^{\sigma t^*} \sin \pi z^*.$$

Substituting Eqs (4.2) in linearized perturbation dimensionless Eqs (3.1)-(3.4) and dropping asterisks for convenience, we get the following equations

$$\begin{aligned}
& \left\{ \left(\sigma + \left((1 - Vz^2)/k \right) \right) (D^2 - a^2)^2 w_1(z) - (1 - Vz^2) (D^2 - a^2)^2 w_1(z) \right. \\
& \left. + 2V (D^2 - a^2) [I + 2z D w_1(z)] - \frac{1}{k} 2vz D w_1(z) \right\} A + a R^{1/2} [I + M_1 (I + M_5) (I - S_T) T_1(z)] B + \\
& - a R_S^{1/2} (I + M_4 + M_4 M_5^{-1}) S_1(z) C + a R^{1/2} M_1 (I + M_5) D \phi_1(z) F = 0,
\end{aligned} \quad (4.3)$$

$$aR^{1/2}(1-M_2-M_2M_5)w_1(z)A+(D^2-a^2-P_r\sigma)T_1(z)B+P_r\sigma M_2D\phi_1(z)F=0, \quad (4.4)$$

$$\begin{aligned} & -aR_S^{1/2}M_6w_1(z)A+S_T M_5 M_6^{-1}(R/R_S)^{1/2}(D^2-a^2)T_1(z)B+ \\ & +\left[\tau(D^2-a^2)-\sigma P_r\right]S_1(z)C=0, \end{aligned} \quad (4.5)$$

$$-R_S^{1/2}\pi^2(1-S_T)DT_1(z)B+R^{1/2}M_5M_6^{-1}DS_1(z)C+R_S^{1/2}(D^2-a^2M_3)\phi_1(z)F=0, \quad (4.6)$$

For the existence of non-trivial solutions, the determinant of the coefficients of A , B , C and F must vanish. This determinant on simplification yields

$$-T_1\sigma^3+T_2\sigma^2+T_3\sigma+T_4=0 \quad (4.7)$$

where

$$T_1=\langle P_r S_1 \langle P_r T_1 \ b_1 w_1 \rangle \ b_5 \phi_1 \rangle,$$

$$T_2=\langle P_r T_1 \langle \tau b_1 S_1 \ b_1 w_1 \rangle \ b_5 \phi_1 \rangle - \langle b_1 T_1 \langle P_r S_1 \ b_1 w_1 \rangle \ b_5 \phi_1 \rangle,$$

$$\begin{aligned} T_3 = & -\langle b_1 T_1 \langle \tau b_1 S_1 \ b_1 w_1 \rangle \ b_5 \phi_1 \rangle + \left\langle \left(1-Vz^2\right) \frac{b_1}{k} w_1 \langle P_r T_1 \ \tau b_1 S_1 \rangle \ b_5 \phi_1 \right\rangle + \\ & -\left\langle \left(1-Vz^2\right) b_1^2 w_1 \langle P_r T_1 \ \tau b_1 S_1 \rangle \ b_5 \phi_1 \right\rangle + \langle 2Vb_1(1+2zDw_1) \langle P_r T_1 \ \tau b_1 S_1 \rangle \ b_5 \phi_1 \rangle + \\ & -\left\langle \frac{2Vz}{k} Dw_1 \langle P_r T_1 \ \tau b_1 S_1 \rangle \ b_5 \phi_1 \right\rangle + a^2 R \langle b_2 D\phi_1 \langle P_r T_1 \ M_6 w_1 \rangle \ b_4 DS_1 \rangle + \\ & -\left\langle \left(1-Vz^2\right) \frac{b_1}{k} w_1 \langle b_1 T_1 \ P_r S_1 \rangle \ b_5 \phi_1 \right\rangle + a^2 R_S \langle b_3 S_1 \langle P_r T_1 \ M_6 w_1 \rangle \ b_5 \phi_1 \rangle + \\ & +\left\langle \left(1-Vz^2\right) b_1^2 w_1 \langle b_1 T_1 \ P_r S_1 \rangle \ b_5 \phi_1 \right\rangle - \langle 2Vb_1(1+2zDw_1) \langle b_1 T_1 \ P_r S_1 \rangle \ b_5 \phi_1 \rangle + \\ & +\left\langle \frac{2Vz}{k} Dw_1 \langle b_1 T_1 \ P_r S_1 \rangle \ b_5 \phi_1 \right\rangle + a^2 R \langle (1+b_2(1-S_T)T_1) \langle w_1 \ P_r S_1 \rangle \ b_5 \phi_1 \rangle + \\ & +a^2 R \langle b_2 D\phi_1 \langle w_1(1-S_T)DT_1 \rangle P_r S_1 \rangle, \end{aligned}$$

$$\begin{aligned} T_4 = & -\left\langle \left(1-Vz^2\right) \frac{b_1}{k} w_1 \langle b_1 T_1 \ \tau b_1 S_1 \rangle \ b_5 \phi_1 \right\rangle + \left\langle \left(1-Vz^2\right) b_1^2 w_1 \langle b_1 T_1 \ \tau b_1 S_1 \rangle \ b_5 \phi_1 \right\rangle + \\ & -\langle 2Vb_1(1+2zDw_1) \langle b_1 T_1 \ \tau b_1 S_1 \rangle \ b_5 \phi_1 \rangle + \left\langle \frac{2Vz}{k} Dw_1 \langle b_1 T_1 \ b_5 \phi_1 \rangle \tau b_1 S_1 \right\rangle + \\ & +a^2 R \langle (1+b_2(1-S_T)T_1) \langle w_1 \ \tau b_1 S_1 \rangle \ b_5 \phi_1 \rangle + a^2 R_S \langle b_3 S_1 \langle w_1 \ S_T b_1 b_4 T_1 \rangle \ b_5 \phi_1 \rangle + \\ & +a^2 R_S \langle b_3 S_1 \langle P_r T_1 \ M_6 w_1 \rangle \ b_5 \phi_1 \rangle + a^2 R \langle b_2 D\phi_1 \langle w_1 \ S_T b_1 b_4 T_1 \rangle \ b_4 DS_1 \rangle + \\ & +a^2 R \langle b_2 D\phi_1 \langle w_1(1-S_T)DT_1 \rangle \tau b_1 S_1 \rangle + a^2 R \langle b_2 D\phi_1 \langle b_1 T_1 \ M_6 w_1 \rangle \ b_4 DS_1 \rangle, \end{aligned}$$

$$b_1 = D^2 - a^2, \quad b_2 = M_1(I + M_5), \quad b_3 = I + M_4 + (M_4/M_5),$$

$$b_4 = (M_5/M_6) \quad \text{and} \quad b_5 = D^2 - a^2 M_3.$$

For obtaining stationary instability, the time-dependent term T_4 is equal to zero. From Eq.(4.7) it is easy to obtain the eigen value R_c .

$$R_c = \frac{x_1 - a^2 R_S (x_2 S_T + x_3)}{\tau x_4 + M_1 (I + M_5) [(I - S_T) \tau x_5 + x_6 S_T + \tau (I - S_T) x_7 + x_8]}$$

where

$$x_1 = -\left\langle (1 - Vz^2) \frac{b_1}{k} w_1 \langle b_1 T_1 \quad \tau b_1 S_1 \rangle b_5 \phi_1 \right\rangle + \left\langle (1 - Vz^2) b_1^2 w_1 \langle b_1 T_1 \quad \tau b_1 S_1 \rangle b_5 \phi_1 \right\rangle + \\ - \left\langle 2Vb_1 (1 + 2zDw_1) \langle b_1 T_1 \quad \tau b_1 S_1 \rangle b_5 \phi_1 \right\rangle + \left\langle \frac{2Vz}{k} Dw_1 \langle b_1 T \quad b_5 \phi_1 \rangle \tau b_1 S_1 \right\rangle,$$

$$x_2 = \langle b_3 S_1 \langle w_1 \quad b_1 b_4 T_1 \rangle b_5 \phi_1 \rangle,$$

$$x_3 = \langle b_3 S_1 \langle P_r T_1 \quad M_6 w_1 \rangle b_5 \phi_1 \rangle,$$

$$x_4 = \langle I \langle w_1 \quad \tau b_1 S_1 \rangle b_5 \phi_1 \rangle,$$

$$x_5 = \langle T_1 \langle w_1 \quad \tau b_1 S_1 \rangle b_5 \phi_1 \rangle,$$

$$x_6 = \langle D\phi_1 \langle w_1 \quad b_1 b_4 T_1 \rangle b_4 DS_1 \rangle,$$

$$x_7 = \langle D\phi_1 \langle w_1 \quad DT_1 \rangle b_1 S_1 \rangle,$$

$$x_8 = \langle D\phi_1 \langle b_1 T_1 \quad M_6 w_1 \rangle b_4 DS_1 \rangle$$

where $\langle u, v \rangle = \int_{-1/2}^{1/2} uv dz$ and w_1, T_1, ϕ_1 and S_1 are trial functions that satisfy the boundary conditions. The

above choice of trigonometry function tacitly implies the use of a higher order Galerkin method. For very large M_1 , one gets the results for the magnetic mechanism, and the critical thermomagnetic Rayleigh number for stationary mode is calculated using

$$N_c = M_1 R_c = \frac{x_1 - a^2 R_S (x_2 S_T + x_3)}{(I + M_5) [(I - S_T) \tau (x_5 + x_7) + x_6 S_T + x_8]}.$$

5. Discussion of results

The linear stability analysis of Soret driven thermohaline convection in a ferromagnetic fluid layer heated from below and salted from above saturating a porous medium subjected to a transverse uniform magnetic field has been considered in the presence of temperature dependent viscosity by using the Brinkman model. Here the free-free boundary conditions are used. The present investigation is carried out through stationary instability. The small thermal perturbation technique is used and the normal mode technique is applied for the perturbation quantities.

Before we discuss the significant results of the system, we turn our attention to the possible range of values of various parameters arising in the study. The range of values of the temperature dependent viscosity parameter V is assumed from 0.1 to 0.5 (Ramanathan and Muchikel, 2006). The ratio of magnitude to gravitational force M_1 , is assumed to be 1000 (Finlayson, 1970). The range of salinity Rayleigh number R_S is between -500 and 500 and Soret parameter S_T ranges from -0.002 to 0.002 (Sekar *et al.*, 2013). The Brinkman model has been used for the permeability k which ranges from 0.1 to 0.9 (Vaidyanathan *et al.*, 2005) and the non-buoyancy magnetization parameter M_3 is taken from 5 to 25 (Sekar *et al.*, 2013). For these type of fluids M_2 will have a negligible value and hence taken to be zero. The Prandtl number P_T is taken to be 0.01 (Vaidyanathan *et al.*, 2005) and the magnetic numbers M_4 , M_5 and M_6 are assumed to be 0.1 (Sekar *et al.*, 2013 and Sekar and Raju, 2013). The ratio of mass transport to heat transport τ is assumed from 0.03 to 0.011 (Sekar *et al.*, 2006).

Figure 1 presents the plots of the critical magnetic thermal Rayleigh number N_C versus the non-buoyancy magnetization parameter M_3 for different values of the temperature dependent viscosity V , $S_T = -0.002$, $R_S = -500$ and $k = 0.1$. It indicates that the non-buoyancy magnetization parameter M_3 has a destabilizing effect on the system when both V and M_3 are increased. This is shown by a fall in N_C values. This is because variation in magnetization releases extra energy which adds up to the thermal energy to destabilize the system.

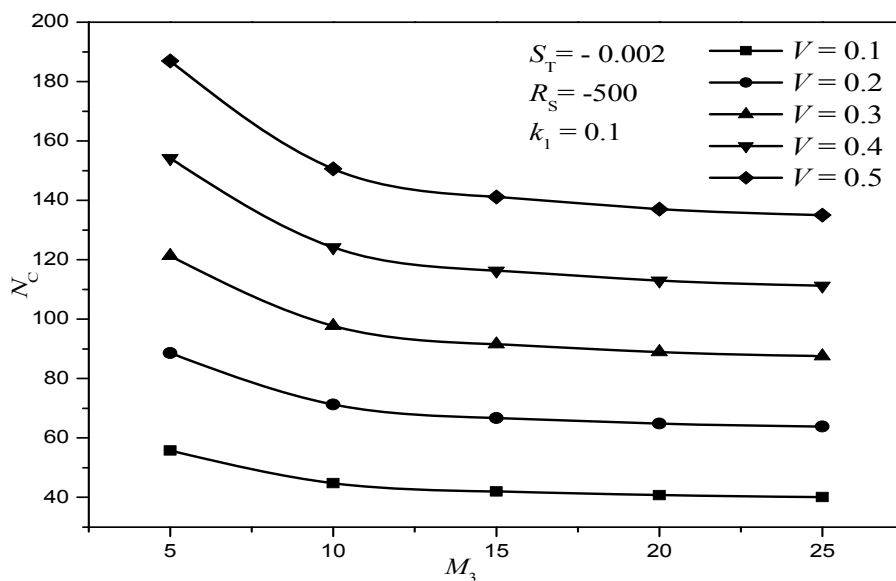
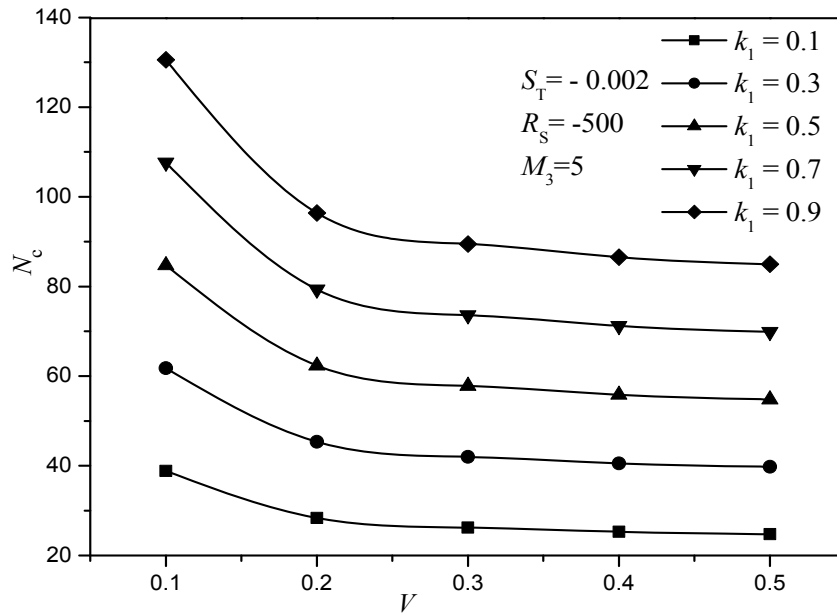


Fig.1. Marginal instability curve for variation of N_C versus M_3 for different values of the temperature dependent viscosity V .

In Figs 2a and b, the variation of the critical thermal magnetic Rayleigh number N_C versus the temperature dependent viscosity V for different permeability of the porous medium k . Both figures exhibit a destabilizing behavior because the presence of a porous medium increases from 0.1 to 0.9 , N_C decreases. It is also observed from the figures that the increase in the pore size makes the fluid flow easy to cause

convection early. Figure 2a illustrates that as k increases, N_C decreases for a negative range of R_S and S_T and this behavior can also be observed exactly in the positive range of R_S and S_T in Fig.2b. Therefore, Figs 2a and b illustrate the same destabilizing effect on the convective system.

(a)



(b)

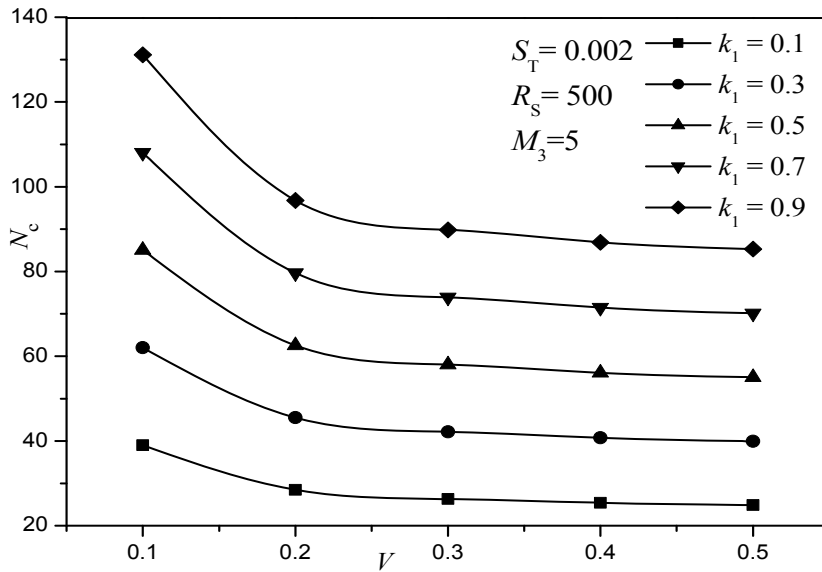


Fig.2. (a) Marginal instability curve for variation of N_C versus V for different values of the medium permeability k , $S_T = -0.002$, $R_S = -500$, $\tau = 0.03$, $M_3 = 5$. (b) Marginal instability curve for variation of N_C versus V for different values of the medium permeability k , $S_T = 0.002$, $R_S = 500$, $\tau = 0.03$, $M_3 = 5$.

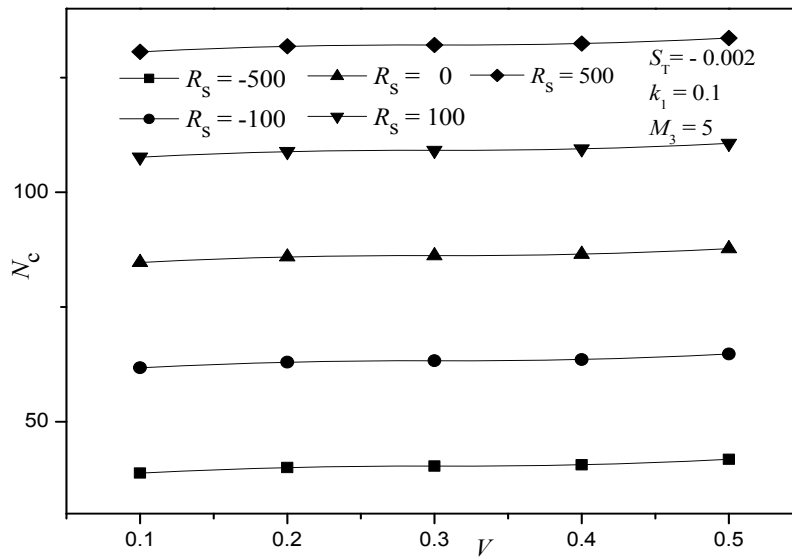
Figures 3a, b and c represent N_C versus V for different values of R_S , S_T and k . It is observed from Figs 3a and b that the temperature dependent viscosity V has a stabilizing effect on the system when V increases, N_C increases and this stabilizing effect of V is rather pronounced. Further, both Figs 3a and b are analyzed for different porous medium k . When the values of k are 0.1 and 0.9 and negative value of S_T , the system show

the same stabilizing behavior and there is no change in the convective system. Figure 3c illustrates N_C versus V for the positive value of S_T , due to the positive value of S_T and $k = 0.9$, the system has a non-equilibrium position compared with Figs 3a and b. Also in Fig.3c, the destabilizing behavior is not much pronounced when V increases, N_C decreases.

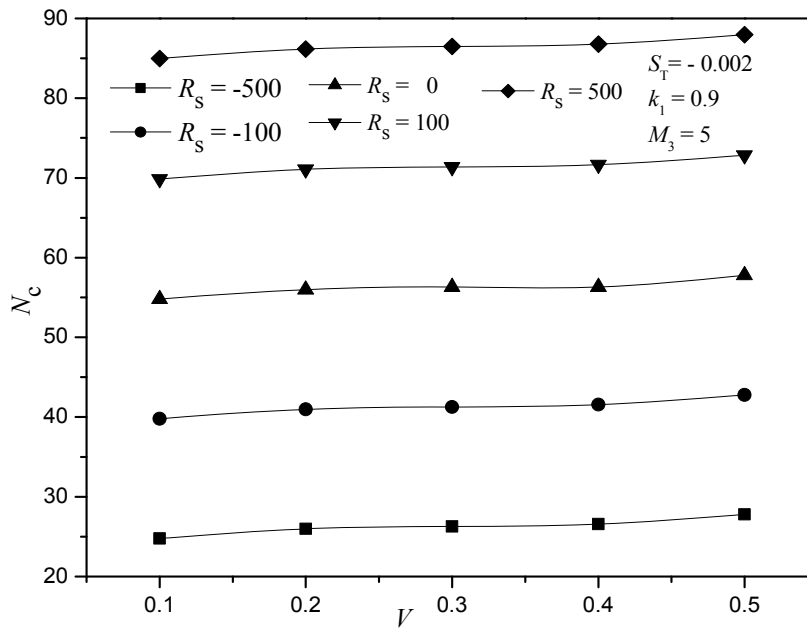
Figures 4a, b and c show the variation of N_C versus the interdiffusion of heat and mass, namely the Soret effect S_T for different V , R_S and k . Figures 4a and b give as increase of S_T , increase of N_C . This leads to stabilizing effect is not much pronounced. Figure 4c show that as S_T increases, N_C decreases. It is seen that the system destabilizes.

It is observed from Fig.5 that the increase in the ratio of mass transport to the heat transport τ shows a stabilizing behavior, for an increasing value of V . When positive range of R_S and S_T , the critical magnetic Rayleigh number N_C has the equal value of the negative range of R_S and S_T . This is because the increase in mass transport adds up to the system to be top heavy.

(a)



(b)



(c)

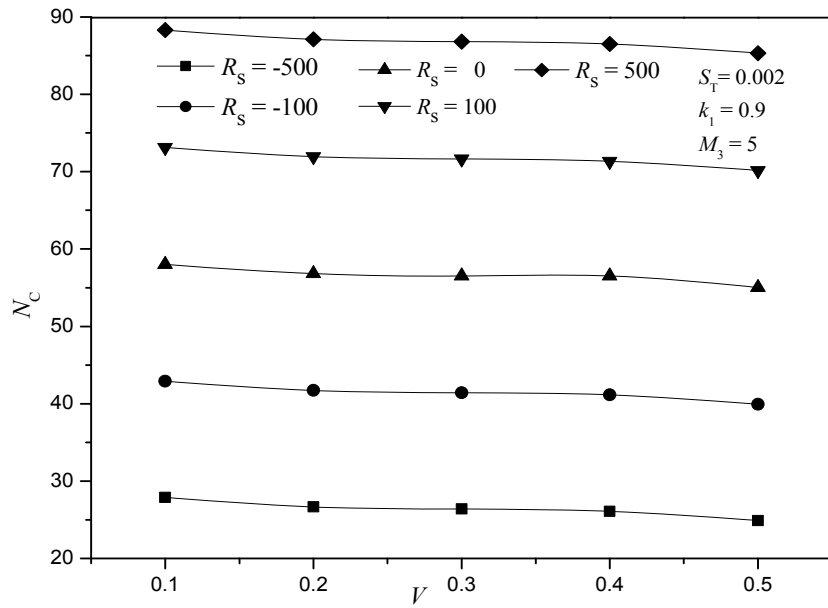
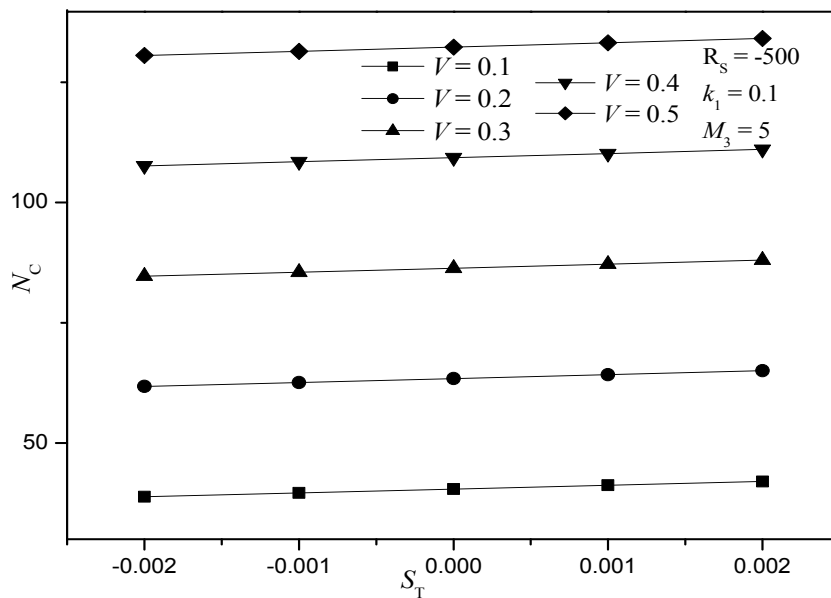
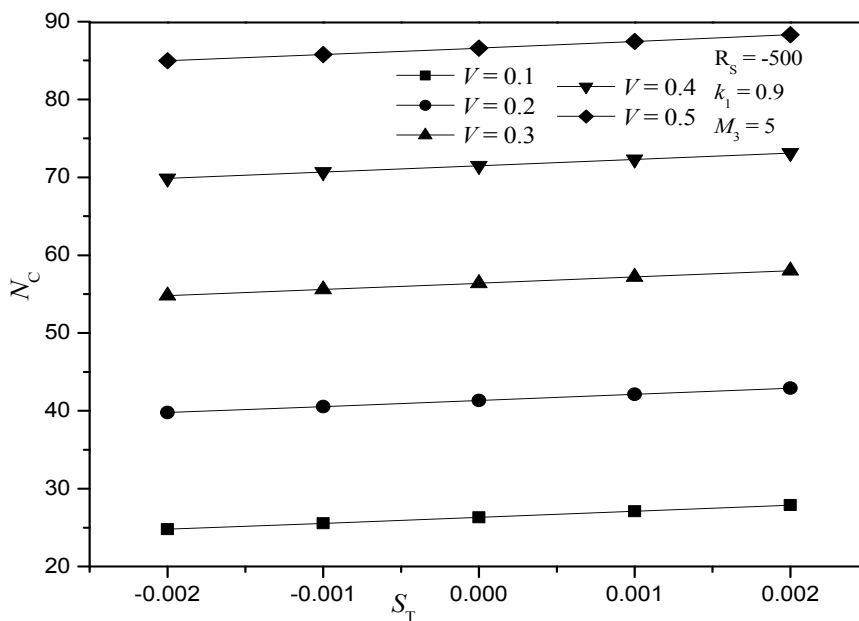


Fig.3. (a) Marginal instability curve for variation of N_C versus V for different values of the salinity Rayleigh number R_S , $S_T = -0.002$, $M_3=5$, $\tau = 0.03$ and $k = 0.1$. (b) Marginal instability curve for variation of N_C versus V for different values of the salinity Rayleigh number R_S , $S_T = -0.002$, $M_3=5$, $\tau = 0.03$ and $k = 0.9$. (c) Marginal instability curve for variation of N_C versus V for different values of the salinity Rayleigh number R_S , $S_T = 0.002$, $M_3=5$, $\tau = 0.03$ and $k = 0.9$.

(a)



(b)



(c)

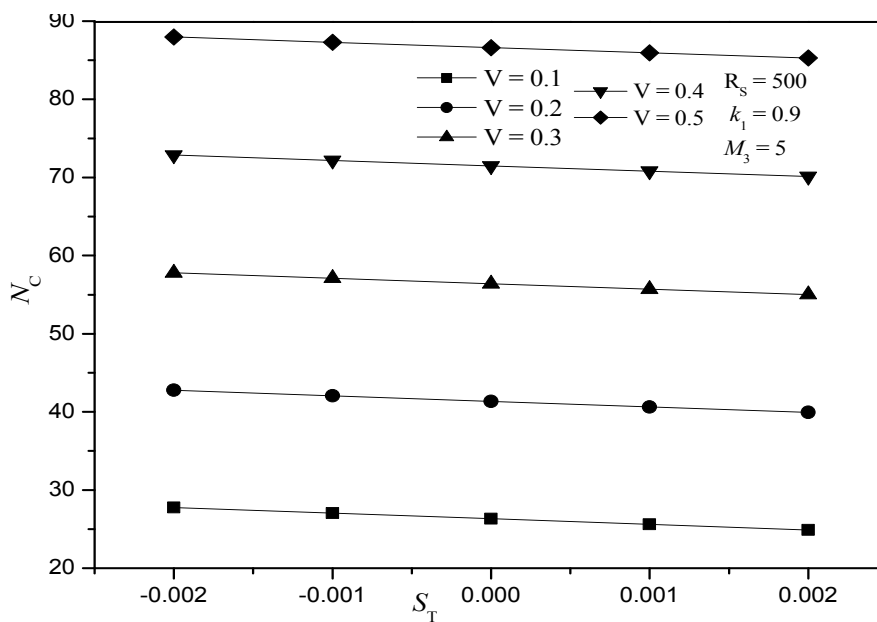


Fig.4. (a) Marginal instability curve for variation of N_C versus S_T for different values of the temperature dependent viscosity V , $R_S = -500$, $M_3 = 5$, $\tau = 0.03$ and $k = 0.1$. (b) Marginal instability curve for variation of N_C versus S_T for different values of the temperature dependent viscosity V , $R_S = -500$, $M_3 = 5$, $\tau = 0.03$ and $k = 0.9$. (c) Marginal instability curve for variation of N_C versus S_T for different values of the temperature dependent viscosity V , $R_S = 500$, $M_3 = 5$, $\tau = 0.03$ and $k = 0.9$.

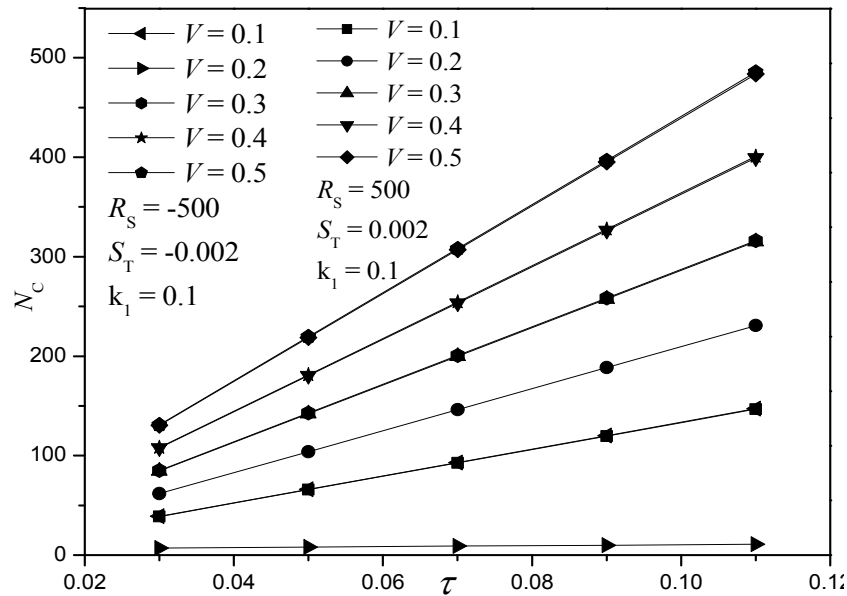


Fig.5. Marginal instability curve for variation of N_C versus τ for different values of the temperature dependent viscosity V , $M_3=5$, $k=0.1$, $R_S=-500$, $S_T=-0.002$ and $R_S=500$, $S_T=0.002$.

6. Conclusion

Soret-driven ferro thermoconvective instability of a magnetic fluid layer heated from below and salted from above in the presence of a porous medium and temperature field dependent viscosity suspended to a transverse uniform magnetic field has been investigated using the Brinkman model. The computational Galerkin method is applied. In this investigation, we have analyzed the effect of various parameters such as the medium permeability, buoyancy magnetization parameter, non-buoyancy magnetization parameter, ratio of mass transport to heat transport, temperature dependent viscosity parameter, Soret coefficient, salinity Rayleigh number and wave number.

The non-buoyancy magnetization parameter M_3 and the permeability of the porous medium k have a destabilizing influence on the convective system. The stabilizing effect is investigated for the temperature dependent viscosity parameter V in a very small value of salinity concentration and also the destabilizing behavior is analyzed for the temperature dependent viscosity parameter V in a large value of salinity concentration.

Thus from the above analysis, one can conclude that the magnetization parameter, temperature gradient and salinity gradient have a profound influence on the onset of convection in a porous medium.

Acknowledgement

The authors are grateful to Prof. D. Govindarajulu, Principal, Pondicherry Engineering College, Puducherry for his constant encouragement. The author K. Raju is thankful to UGC for grant of Rajiv Gandhi National Fellowship 2010 – 2011 (Award letter number: F. 14-2(SC)/2010 (SA-III), Dated: May 2011).

Nomenclature

- B – magnetic induction T
 $C_{v,H}$ – effective heat capacity at constant volume and magnetic field (kJ/m^3K)

- D/Dt – convective derivative $s^{-1} [D/Dt = \partial/\partial t + \mathbf{q} \cdot \nabla]$
- d – thickness of the fluid layer m
- \mathbf{g} – gravitational acceleration $(0, 0, -g) ms^{-2}$
- \mathbf{H} – magnetic field amp/m
- K – mass diffusivity
- K – pyromagnetic coefficient $\left[\equiv -(\partial M/\partial T)_{H_0, T_0} \right]$
- K_l – thermal diffusivity W/mK
- K_2 – salinity magnetic coefficient $\left[\equiv -(\partial M/\partial S)_{H_0, T_0} \right]$
- K_s – concentration diffusivity W/mkg
- k – permeability of the porous medium
- k_0 – resultant wave number $\left[k_0 = \sqrt{k_x^2 + k_y^2} \right] m^{-1}$
- k_x, k_y – wave number in the x and y direction m^{-1}
- \mathbf{M} – magnetization $Ampm^{-1}$
- M_0 – mean value of the magnetization at $H = H_0$ and $T = T_0$
- P – hydrodynamic pressure (N/m^2)
- \mathbf{q} – velocity of the ferrofluid $(u, v, w) ms^{-1}$
- S – solute concentration kg
- S_T – Soret coefficient
- T – temperature K
- t – time s
- α_t – coefficient of thermal expansion K^{-1}
- α_s – analogous solvent coefficient of expansion K^{-1}
- β_t – uniform temperature gradient Km^{-1}
- β_s – uniform concentration gradient kgm^{-1}
- μ_0 – magnetic permeability of vacuum
- μ_l – reference viscosity at $T = T_0$
- μ – dynamic viscosity $kgm^{-1}s^{-2}$
- ρ_0 – mean density of the clean fluid kgm^{-3}
- ρ – density of the fluid kgm^{-3}
- σ – growth rate s^{-1}
- φ – viscous dissipation factor containing second order terms in velocity
- ϕ – magnetic scalar potential Amp
- θ – perturbation in temperature (K)
- χ – magnetic susceptibility $\left[\equiv -(\partial M/\partial H)_{H_0, T_0} \right]$
- ∇ – vector different operator $\left[\equiv \mathbf{i}(\partial/\partial x) + \mathbf{j}(\partial/\partial y) + \mathbf{k}(\partial/\partial z) \right]$

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Received: August 13, 2013

Revised: December 15, 2013